

A highly accurate ab initio dipole moment surface for the ground electronic state of water vapour for spectra extending into the ultraviolet

Eamon K. Conway, Aleksandra A. Kyuberis, Oleg L. Polyansky, Jonathan Tennyson, and Nikolai F. Zobov

Citation: *The Journal of Chemical Physics* **149**, 084307 (2018); doi: 10.1063/1.5043545

View online: <https://doi.org/10.1063/1.5043545>

View Table of Contents: <http://aip.scitation.org/toc/jcp/149/8>

Published by the [American Institute of Physics](#)

PHYSICS TODAY

WHITEPAPERS

ADVANCED LIGHT CURE ADHESIVES

Take a closer look at what these environmentally friendly adhesive systems can do

READ NOW

PRESENTED BY
 **MASTERBOND**
ADHESIVES | SEALANTS | COATINGS

A highly accurate *ab initio* dipole moment surface for the ground electronic state of water vapour for spectra extending into the ultraviolet

Eamon K. Conway,^{1,2} Aleksandra A. Kyuberis,³ Oleg L. Polyansky,^{1,3} Jonathan Tennyson,¹ and Nikolai F. Zobov³

¹Department of Physics and Astronomy, University College London, Gower Street, London WC1E 6BT, United Kingdom

²Atomic and Molecular Physics Division, Harvard-Smithsonian Center for Astrophysics, Cambridge, Massachusetts 02138, USA

³Institute of Applied Physics, Russian Academy of Sciences, Uljanov Street 46, Nizhniy Novgorod 603950, Russia

(Received 10 June 2018; accepted 8 August 2018; published online 27 August 2018)

A new global and highly accurate *ab initio* dipole moment surface (DMS) for water vapour is presented. This DMS is based on a set of 17 628 multi-reference configuration interaction data points that were calculated with the aug-cc-pCV6Z basis set with the Douglas-Kroll-Hess Hamiltonian; tests are performed at several other levels of *ab initio* theory. This new “CKAPTEN” DMS improves agreement with recent experimental measurements compared with previous models that poorly predicted some bands in the infrared while also maintaining or improving on the agreement for all remaining strong lines. For high overtones located in both the visible and the near ultraviolet regions, our predicted intensities all lie within 10% of recent atmospheric observations. A crossing of energy levels in the ν_1 fundamental and $2\nu_2$ states is seen to offset transition intensities in the ν_1 fundamental band; residual inaccuracies within the potential energy surface used is the cause of this problem. © 2018 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). <https://doi.org/10.1063/1.5043545>

I. INTRODUCTION

Water vapour can certainly be regarded as the most important molecule on Earth; it is one of the basic requirements for life to exist¹ but is also the major contributor to the greenhouse effect and the largest absorber of incoming sunlight.² It has long been suggested that climate change may have significant implications on our hydrological cycle,³ either suppressing it⁴ or possibly enhancing it.⁵ Similarly, any changes in atmospheric water vapour concentration will affect our climate.

The key to understanding this complex relationship focuses on our ability to both accurately measure and model H₂¹⁶O spectra and hence water vapour concentrations. Atmospheric observations are usually based on spectroscopic databases such as HITRAN⁶ or GEISA.⁷ Given the abundance of water vapour and the complexity surrounding its spectrum, water absorption lines can often obscure spectral features due to trace molecules whose monitoring can be important for a variety of atmospheric issues.

The upcoming NASA Tropospheric Emissions: Monitoring of Pollution (TEMPO) mission⁸ plus the related European Sentinel⁹ and Korean GEMS (geostationary environmental monitoring satellite)¹⁰ missions will probe the Earth's atmosphere in the blue visible and near ultraviolet. As Lampel and co-workers^{11,12} have already demonstrated, retrievals for trace species in this region require dramatically improved understanding of the underlying water absorptions. Furthermore, it has been argued that retrieval of water columns at these

short wavelengths has significant advantages^{13,14} compared with measurements using longer wavelengths. This argument may have consequences for previous satellite missions that had scanned the Earth's atmosphere in the visible and near UV, such as GOME,¹⁵ SCIAMACHY,¹⁶ GOME-2,¹⁷ and OMI.¹⁸ However, precise retrievals rely on the availability of accurate laboratory data. Even in this region, the absorption by water vapour is due to weak, high overtone rotation-vibration transitions, on which data are largely lacking.

Recently, Birk *et al.*¹⁹ measured new experimental water line intensities in the infrared region and followed this up with a thorough comparison of these new accurate measurement against a computed line list that used variational nuclear motion calculations and the LTP2011 dipole moment surface (DMS) of Lodi *et al.*,²⁰ the current release of HITRAN, HITRAN2016, contains large sections based on transition intensities computed using the LTP2011 DMS.^{21,22} Birk *et al.*'s findings are significant, as while they found excellent agreement for many transitions, they also identified a number of discrepancies including the computed intensities of the ν_1 fundamental symmetric stretching band that deviated from experiment by between +5% and -13%; such disagreements for the fundamental are well known.^{23,24}

Recent atmospheric measurements in the visible and near ultraviolet by Lampel *et al.*¹¹ detected for the first time a prominent H₂¹⁶O absorption band at 363 nm, which corresponds to the eighth overtone of stretch quanta, (9,0)[±]0. Lampel *et al.* showed that neglecting this water absorption had a significant effect on the accurate retrieval of important trace molecules

such as O₄, HONO, SO₂, and OClO, which show absorption features close to those of water vapour. Lampel *et al.* also suggested that water vapour absorption at wavelengths shorter than 360 nm may also affect the retrieval of BrO and HCHO.

Notable among these *ab initio* DMSs are available.^{20,23–28} Notable among these are the SP2000 surface of Schwenke and Partridge,²⁴ which was used in the BT2 line list,²⁹ and a number of surfaces from Lodi *et al.*: Core-Valence-Relativistic (CVR)²⁸ and the more recent Lodi-Tennyson-Polyansky LTP2011 and LTP2011S²⁰ models that provide a full and reduced (smaller) fit to the same *ab initio* dataset.

Lampel *et al.*¹¹ used transition intensity calculations based on several of these DMSs to interpret their retrievals and none of them provided accurate results. The most recent and complete POKAZATEL water line list of Polyansky *et al.*³⁰ used the LTP2011S surface and showed good behavior around 363 nm but was too weak, while the earlier CVR surface²⁸ had an irregular shape but demonstrated stronger absorption than LTP2011S. Lampel *et al.* concluded that the shape of water absorption in the 340–380 nm range is better replicated by POKAZATEL than by CVR, BT2, or HITEMP.³¹

Broadband, laboratory measurements by Du *et al.*³² of near ultraviolet water cross sections show very significant absorption which Wilson *et al.*³³ could not replicate. On the basis of their measurements, Du *et al.* also make assertions about the importance of short-wavelength absorption by water vapour to the Earth's energy budget. The precise value of the absorption by the vibration-rotation spectrum of water in the near ultraviolet remains an important open question. At present, the maximum transition frequency for H₂¹⁶O within HITRAN2016 is 25 710.8 cm⁻¹, in line with high-resolution laboratory studies where intensity measurements stop in this region.^{34,35}

The difficulty in creating an *ab initio* DMS capable of accurately modeling weak absorption in high-energy regions is of no surprise.^{24,36,37} Possible sources of error are extensive and range from the accuracy of the underlying *ab initio* calculations, the number of data points calculated, the choice of functional form used in the fitting procedure, and the number of parameters used in the fit. The purpose of this work is to create a global DMS for the ground electronic state of the water molecule that improves the accuracy of previous models while simultaneously solving the issues surrounding high-frequency transitions occurring in both the visible and the ultraviolet.

II. COMPUTATIONAL METHODS

All electronic structure computations were carried out with the quantum chemistry package MOLPRO.³⁸ Dipoles were computed using the finite differences (FD) approach, which necessitated four calculations per point for the dipoles and one other at zero field to obtain the energy at that geometry. While being computationally more demanding, experience has shown us that this method does, however, provide more accurate dipoles than the alternative scenario of taking an expectation value (XP).³⁹ It also facilitates the inclusion of the contribution to the dipole of small corrections

whose contribution to the energy is treated using perturbation theory.

Dipoles were calculated at the multi-reference configuration interaction (MR-CI) level of theory. To ensure numerical stability in the FD, the energy convergence threshold was set to 5×10^{-11} E_h for all MR-CI calculations. The relaxed reference Davidson correction (+Q) has been applied to the MR-CI dipoles. Unless otherwise stated, relativistic corrections were obtained with the mass-velocity Darwin one-electron (MVD1) operator available within MOLPRO.

Four different data sets were computed: one with the aug-cc-pV5Z basis set, another with the aug-cc-pwCV5Z basis set, and two using the larger aug-cc-pCV6Z basis set.^{40–42} With the exception of valence-only calculations, all ten electrons were otherwise fully correlated with both the complete active space self-consistent field (CAS-SCF) and MR-CI methods using an active space consisting of the lowest nine energetic A' molecular orbitals (MO's) and the two lowest A'' MO's, denoted as (9,2).

A. Dipole accuracy

Lodi *et al.*²⁸ selected 14 “key structure” geometries that overall offer an insight into the global behavior of the calculation methods; these same 14 points were used for our own investigation. Here and elsewhere, r_1 and r_2 identify the bond length between the oxygen atom centred at the origin and both hydrogen atoms, with θ representing \angle HOH. First, we investigated the differences between dipoles computed under two different electric field strengths, 5×10^{-5} a₀ and 3×10^{-4} a₀, for basis sets aug-cc-pCV(T,Q,5,6)Z and aug-cc-pwCV5Z.

Table I shows that the difference between dipoles computed with the two field strength is small and of the order of 10⁻⁶ a.u. Dipoles are, on average, slightly larger when computed with the smaller field strength. The two-point central point difference formula is valid in the limit of the field strength going to zero; hence, the weaker field is preferred if numerical stability is achieved.

1. aug-cc-pwCV5Z

A total of 8921 dipoles have been computed with bond lengths ranging from 1.3141 a₀ to 2.3141 a₀ and $84.35^\circ \leq \theta \leq 124.35^\circ$, of which the maximum energy calculated with respect to equilibrium is 45 789 cm⁻¹. This surface is referred to as pwCV5Z below.

2. aug-cc-pCV6Z

This data set includes 2915 configurations with the length of OH bonds ranging from 1.3891 a₀ to 2.3141 a₀ and $84.35^\circ \leq \theta \leq 124.35^\circ$. The maximum energy within the sample is 23 625 cm⁻¹. This surface is referred to as MVD1 6Z below.

3. aug-cc-pV5Z

For this group, we computed an extended grid, comprising 21 879 molecular configurations whose energies extend to 46 091 cm⁻¹. OH distances stretch from 1.3141 a₀ to 2.41 a₀ with \angle HOH restricted to $50^\circ \leq \theta \leq 178^\circ$. These calculations were performed using a smaller active space, (7,2) with the CAS-SCF requiring a double occupied 1s(O). Based on the result of Lodi *et al.*,²⁸ who found that the (MVD1) relativistic

TABLE I. Differences in dipole moments (in a.u.) computed for the 14 key structure geometries of Lodi *et al.*²⁸ for different electric field strengths using the aug-cc-pCV6Z basis set. Dipole values are shown for the larger field strength of $5 \times 10^{-5} a_0$, with Δ defined as $\mu(5 \times 10^{-5} a_0) - \mu(3 \times 10^{-4} a_0)$.

$r_1(a_0)$	$r_2(a_0)$	θ (deg)	E (cm^{-1})	μ_{\parallel}	μ_{\perp}	$\Delta_{\parallel}(\times 10^{-5})$	$\Delta_{\perp}(\times 10^{-5})$
1.8096	1.8096	104.52	0	0.728297	...	0.5	...
1.7008	1.7008	104.52	1 609.19	0.717352	...	0.0	...
2.3622	2.3622	104.52	19272.34	0.734124	...	0.3	...
1.4173	1.4173	100.00	31534.08	0.703770	...	0.0	...
2.6456	2.6456	100.00	32968.02	0.735414	...	0.0	...
1.7952	1.7952	170.00	10783.94	0.130691	...	0.0	...
1.7952	1.7952	45.00	28064.20	0.967926	...	0.0	...
1.7952	2.3622	179.90	24645.62	0.001 358	-0.330 295	0.0	-0.1
1.6063	2.5511	104.52	18247.82	0.738422	-0.102 659	0.0	0.1
1.7952	3.7795	104.52	37378.02	0.557052	0.188 877	-0.3	0.6
1.8120	1.8120	104.34	0.53	0.729539	...	0.4	...
1.4173	1.6063	120.00	18708.71	0.602082	-0.047 448	0.1	0.0
1.5118	1.7008	100.00	8 959.20	0.731149	-0.028 622	0.0	0.0
2.7132	2.3622	95.00	27424.81	0.780010	-0.011 927	-0.4	0.0

and core corrections almost perfectly cancel, our valence-only dipoles exclude any relativistic treatment. The speed of these computations allowed us to reduce the threshold of MR-CI energy convergence to $2 \times 10^{-11} E_h$, which provides a test of numerical stability of our procedure. This surface is referred to as V5Z below.

4. MVD1 and Douglas-Kroll-Hess comparisons

We investigate differences in dipoles computed with the Douglas-Kroll-Hess Hamiltonian to order two (DKH2) and the MVD1 method for the 14 “key structure” geometries; see Table II. DKH2 dipoles are always larger than the respective MVD1 points, with deviation between the two methods becoming increasingly larger for those molecular configurations whereby one or both bonds are stretched further beyond equilibrium.

These findings are rather unexpected; we cannot identify a previous study that investigated both methods of computing dipole moments in depth. We therefore computed dipoles with both computational methods for two different data sets: Set 1 with one bond fixed at $1.31 a_0$ and the angular separation constant at 104.85° and Set 2 with one bond remaining at $1.6 a_0$ and the angle set at 40° . Energies for Set 1 range from $29\,633 \text{ cm}^{-1}$ to $40\,955 \text{ cm}^{-1}$, with those from Set 2 in the region of $36\,727 \text{ cm}^{-1}$ to $43\,373 \text{ cm}^{-1}$. Figure 1 plots the difference between the two methods for these two sets. The DKH2 dipoles are consistently larger than those from MVD1.

A decreasing trend in deviation is observed at $2.6 a_0$, which corresponds to an approximate molecular configuration energy of $38\,100 \text{ cm}^{-1}$. The reason for such behavior is presently unknown.

TABLE II. Differences between DKH2 and MVD1 dipoles (a.u.) for basis sets 6Z and 5Z calculated by subtracting the MVD1 value from that of DKH2. Shown are the DKH2 dipoles for the aug-cc-pCV6Z basis set.

r_1	r_2	θ	E^a (cm^{-1})	μ_{\parallel}	μ_{\perp}	6Z (10^{-5})		5Z (10^{-5})	
						Δ_{\parallel}	Δ_{\perp}	Δ_{\parallel}	Δ_{\perp}
1.8096	1.8096	104.52	0	0.728311	...	1.4	...	2.4	...
1.7008	1.7008	104.52	1 609.19	0.717367	...	1.5	...	1.7	...
2.3622	2.3622	104.52	19272.34	0.734182	...	5.8	...	6.5	...
1.4173	1.4173	100.00	31534.08	0.703779	...	0.9	...	1.4	...
2.6456	2.6456	100.00	32968.02	0.735502	...	8.8	...	9.9	...
1.7952	1.7952	170.00	10783.94	0.130691	...	0.2	...	0.2	...
1.7952	1.7952	45.00	28064.20	0.967963	...	3.7	...	4.5	...
1.7952	2.3622	179.90	24645.62	0.001 357	-0.330 286	0.1	0.9	0.0	1.0
1.6063	2.5511	104.52	18247.82	0.738467	-0.102 694	4.5	-3.5	5.3	-3.7
1.7952	3.7795	104.52	37378.02	0.557016	0.188 954	-3.6	7.7	-2.8	7.6
1.8120	1.8120	104.34	0.53	0.729554	...	1.5	...	2.2	...
1.4173	1.6063	120.00	18708.71	0.602089	-0.047 449	0.7	-0.1	1.1	-0.2
1.5118	1.7008	100.00	8 959.20	0.731162	-0.028 625	1.3	-0.3	1.8	-0.3
2.7132	2.3622	95.00	27424.81	0.780091	-0.011 906	8.1	2.1	9.0	2.1

^aDKH2 energy with respect to equilibrium geometry using the 6Z basis set.

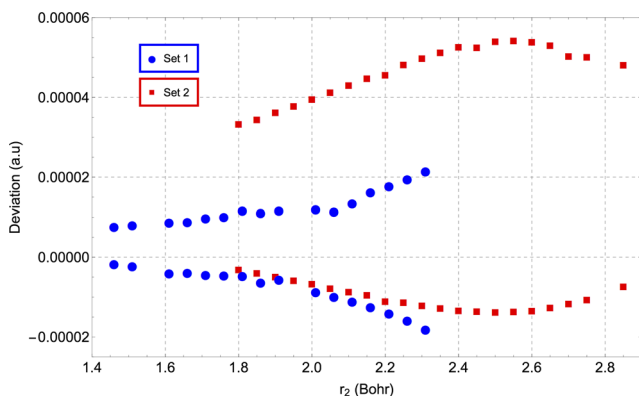


FIG. 1. Differences between parallel and perpendicular dipole components computed with DKH2 and MVD1 methods using the aug-cc-pCV6Z basis set at two different configuration sets. The difference is given by the DKH2 dipole minus the MVD1 dipole. Positive points correspond to parallel dipole components, while negative ones are perpendicular components. Set 1 is for fixed $r_1 = 1.31 a_0$ and $\theta = 104.85^\circ$, while Set 2 has $r_1 = 1.6 a_0$ and $\theta = 40^\circ$.

5. DKH2 aug-cc-pCV6Z

It is also well known that accurate determination of a DMS requires a large number of grid points,^{24,43} hence, the *ab initio* calculations presented here implement the DKH2 method and encompass geometries with $30^\circ \leq \theta \leq 178^\circ$ for a total of 17 628 configurations with OH stretches between $1.3 a_0$ and $4 a_0$. The maximum energy considered was $45\,000 \text{ cm}^{-1}$, and for the linear least squares (LSQ) fitting procedure those points whose energy falls below $32\,000 \text{ cm}^{-1}$ were weighted unity, with the remaining taking a value of 1×10^{-6} . This energy was chosen because points lying below this in energy threshold were not distorted by the introduction of this weighting scheme, while the 3437 *ab initio* points that lie above this threshold help to provide good stability and to constrain the long-range behavior. This surface is referred to as CKAPTEN below.

In total, all *ab initio* calculations accumulate over 110 years of central processing unit (CPU) run time. These were completed in OpenMP and MPI (message passing interface) arrangements on supercomputers *Legion* and *Grace*, respectively, each of which form part of University College London’s (UCL) research computing network. On the *Legion* cluster, Dell C6220 nodes were chosen (hardware subject to small differences as nodes were purchased in sections), of which 6 cores and 90 GB of local disk space were required per data point and each core provided 4 GB of memory. For *Grace*, 2 nodes were utilised, with each node configured with 2×8 core Intel Xeon E5-2630 v3 processors, with 120 GB of SSD (solid-state device) disk space each, and all nodes connected through Intel TrueScale QDR Infiniband infrastructure.

The CKAPTEN surface alone required approximately 80 CPU years of the 110.

III. DMS FIT

The development of a functional form that can accurately model the DMS over an extended range of geometries was not straightforward; it was only via a lengthy process of trial and error that we found a suitable expression.

In the equilibrium configuration, we take $r_e = 1.8141 a_0$ and $\theta_e = 104.52^\circ$, which is close to the experimental value.⁴⁴ The parallel dipole component is denoted as μ_x and bisects the angle θ , and the remaining dipole μ_y is placed perpendicular to μ_x .

Three variables are introduced: $\zeta_1 = \frac{(r_1+r_2)}{2} - r_e$, $\zeta_2 = (r_2 - r_1)$, and $\zeta_3 = \theta/\theta_e$. Each is chosen to approximately represent the behavior of a vibrational mode: ζ_1 reflects the symmetric stretching mode ν_1 , ζ_3 reflects the bending mode ν_2 , and ζ_2 reflects the anti-symmetric stretching mode ν_3 . The analytic expressions used to represent each of the dipole components are as follows:

$$\mu_x(r_1, r_2, \theta) = (\pi - \theta) \sum_{i,j,k} C_{ijk}^{(x)} \zeta_1^i \zeta_2^j \zeta_3^k, \quad (1)$$

$$\mu_y(r_1, r_2, \theta) = \sum_{i,j,k} C_{ijk}^{(y)} \zeta_1^i \zeta_2^j \zeta_3^k. \quad (2)$$

Given our choice of coordinate system, the underlying symmetry provides several constraints which must be adhered to in our functional form. The first requires the parallel dipole to be zero at linear geometries, so we introduce the leading $(\pi - \theta)$ factor, which provides freedom in our choice of the ζ_3 variable. Second, under exchange of r_1 and r_2 , we require μ_x to be symmetric and μ_y to be anti-symmetric. Hence, the exponent of ζ_2 must be even for μ_x and can include zero due to the presence of our pre-leading $(\pi - \theta)$ factor. Similarly for μ_y , the exponent of ζ_2 must be odd.

A total of 126 and 98 parameters were used to fit the parallel and perpendicular dipole components of the aug-cc-pCV6Z CKAPTEN surface, and this gave a weighted root mean square (rms) deviation of 1.5×10^{-4} a.u. and 1.8×10^{-4} a.u., respectively.

In fitting the parallel dipole component of LTP2011,²⁰ Lodi *et al.* obtained a lower residual of 3×10^{-5} a.u. However, while their implementation of 200 fitting parameters to “only” 2628 configurations provides a smaller residual, our data set is over six times larger and we implemented only 128 parameters to ensure stability. For perpendicular dipoles, we obtain a lower rms than the 4×10^{-4} a.u. reported by LTP2011.²⁰

We have successfully restricted the residuals to be less than 10^{-3} a.u. for configurations lying below approximately $30\,000 \text{ cm}^{-1}$ (Fig. 2). This was achieved by reducing the weight

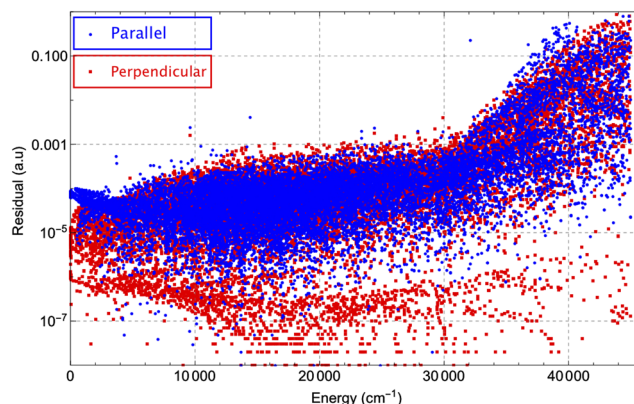


FIG. 2. Log plot of the absolute value of residuals per *ab initio* data point as a function of energy for the CKAPTEN dipole moment surface.

of those geometries whose energy with respect to equilibrium lies above $32\,000\text{ cm}^{-1}$. At $30\,000\text{ cm}^{-1}$, the LTP2011 DMS²⁰ has a deviation of its parallel and perpendicular dipoles of approximately 5×10^{-3} a.u. and 1×10^{-2} a.u., respectively, which would create problems for bands that depend strongly on the perpendicular dipole contribution.

IV. CKAPTEN STABILITY

To test the stability of CKAPTEN DMS to our fitting procedure, we remove half the number of *ab initio* data points and refit this subset using the same functional form. The data points are ordered in increasing energy and every second removed to leave 8814 geometries.

Neither rms value shows any substantial change, with both remaining as they were before to within 1×10^{-5} a.u. Below $10\,000\text{ cm}^{-1}$, we observe less than 1% change in the theoretical intensities of both models.

For known problematic bands¹⁹ above $10\,000\text{ cm}^{-1}$, we note a 1.5% change in the (121) band, with the smaller *ab initio* model predicting stronger line intensities. For the (102) band, we instead see line intensities slightly decrease in magnitude for the smaller model, representing an approximate 1% change. We observe negligible changes in $3\nu_1$ band intensities.

For the atmospheric data of Lampel *et al.*,¹¹ notably the bands at 487 nm (511), 471 nm (303), 377 nm, and 363 nm (900), we see the intensities increase by 5% for the 487 nm (511) band but weaken for the remainders by 8%, 12%, and 11%, respectively. These high energy overtones are extremely difficult to accurately predict, and such changes are small compared with the error associated with previous theoretical models.

V. COMPUTATION OF LINE INTENSITIES

Our line intensities are computed with the well-established, exact kinetic energy operator, variational nuclear motion program DVR3D.⁴⁵

The intensity of a transition $I(\omega_{if})$ in units of cm per molecule, evaluated at frequency ω_{if} , is given by

$$I(\omega_{if}) = \frac{C\omega_{if}g_i[\exp(-E''/kT) - \exp(-E'/kT)]}{Q(T)} \times \left| \sum_{\alpha} \langle i | \mu_{\alpha} | f \rangle \right|^2, \quad (3)$$

where C is a constant of value $4.162\,034 \times 10^{-19}$; E'' and E' are the energy of lower and upper states, respectively; g_i is the total degeneracy factor; and $\langle i | \mu_{\alpha} | f \rangle$ is the expectation value of the α component of the dipole moment between states i and f in Debye. $Q(T)$ is the partition function at temperature T and has been calculated to high accuracy.^{46,47} All intensities quoted in this paper assume the natural isotopologue abundance for H_2^{16}O of 0.997 317, and deviation between theory and experimental line intensities is calculated as $100(\frac{Obs.}{Calc.} - 1)$.

Calculation of the $\langle i | \mu_{\alpha} | f \rangle$ dipole transition matrix elements requires wavefunctions. Unless otherwise stated, these were taken from those generated for the recent

POKAZATEL H_2^{16}O line list.³⁰ The POKAZATEL PES extends the highly accurate potential energy surface (PES) due to Bubukina *et al.*⁴⁸ This provides an accuracy of 0.03 cm^{-1} for energy levels below $25\,000\text{ cm}^{-1}$ and 0.1 cm^{-1} for those between $25\,000\text{ cm}^{-1}$ and $41\,000\text{ cm}^{-1}$.³⁰

We note that Lodi and Tennyson²¹ analyzed the stability of lines with respect to changes in both the DMS and the PES used in the calculations. Stable lines were those that showed little sensitivity and whose uncertainty in their intensity was largely determined by the uncertainty in the underlying DMS. See Zak *et al.*⁴⁹ for a more detailed analysis of this methodology.

Below, where reference is made to intensities of the LTP2011 surface, these lines were computed by us with the POKAZATEL PES.

VI. RESULTS

All intensities are presented for the HITRAN canonical temperature of 296 K. Below, we define weak lines as those with an intensity lower than $10^{-24}\text{ cm molecule}^{-1}$. For transition frequencies in the infrared, we calculate all transitions for $J \leq 20$. Spectral intensities within the visible and ultraviolet regions are confined to $J \leq 14$.

A. Infrared comparisons

Table III documents an average percentage deviation between theoretical intensities of both LTP2011²⁰ and CKAPTEN for the infrared bands measured by both Birk *et al.*¹⁹ and Sironneau and Hodges.⁵⁰ For many bands, the agreement is excellent; below, we focus on those bands that can be regarded as problematic.

Figure 3 compares computed ν_1 line intensities with $J \leq 2$ between various models and shows strong dependence

TABLE III. Band average deviation (%) compared with measured transition intensities of Birk *et al.*¹⁹ and Sironneau and Hodges.⁵⁰ The heading gives the PES used (upper row) and DMS used (lower row). The quoted experimental uncertainty is given for comparison.

Transition	POKAZATEL	POKAZATEL	PES15k	Expt. Unc.
	LTP2011S	CKAPTEN	CKAPTEN	
001-010	-0.42	-0.31	-0.56	0.76
010-000	-0.49	-0.34	-0.26	0.15
100-010	1.24	1.37	1.40	0.83
020-010	-1.70	-1.64	-0.72	0.23
020-000	1.77	1.85	1.29	0.51
030-010	1.72	1.71	2.26	0.29
100-000	2.02	2.10	2.50	0.49
110-010	1.97	2.46	2.11	1.37
011-010	-0.32	1.02	0.70	0.37
001-000	-0.72	0.58	0.66	0.37
101-000	-2.11	-2.63	-2.16	0.32
200-000	0.58	-2.27	-0.92	0.45
002-000	5.25	0.25	0.73	0.22
300-000	-6.05	-1.15	0.25	0.95
121-000	-4.54	-0.84	0.61	0.36
201-000	-1.07	0.01	0.43	0.73
102-000	9.53	1.96	2.94	0.95
003-000	-1.06	2.52	4.15	1.15

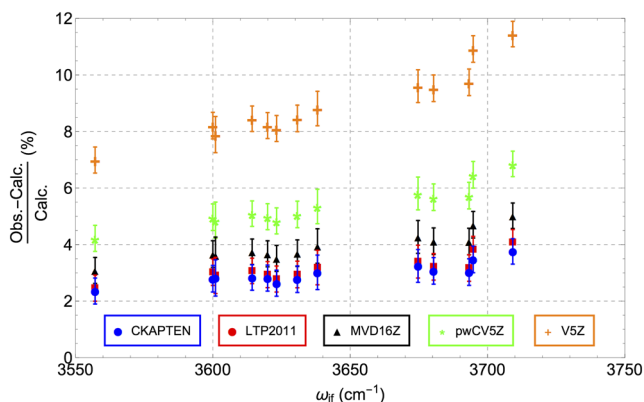


FIG. 3. Intensity deviation from the experimental data of Birk *et al.*¹⁹ in the ν_1 fundamental for transitions with $J \leq 2$ for different theoretical models, LTP2011 of Lodi *et al.*,²⁰ and our best for each basis set.

on the method used to compute the *ab initio* DMS. This is unusual behavior given that these are fundamental transitions and different levels of *ab initio* calculations should not affect this band so dramatically. However, we do note that previous studies have already highlighted this fundamental as being problematic.^{23,24}

The vertical shift between each of the different models is attributed to the underlying electronic structure calculations used to construct the DMS. However, despite significant differences between each of the DMSs tested, we also observe an identical pattern to which the PES, common to all, must be inherently responsible.

Figure 4 recasts these data by intensity and shows that there is a clear change in the deviation between theory and experiment for intensities falling slightly below $10^{-24} \text{ cm molecule}^{-1}$. However, we note that LTP2011 agrees well with CKAPTEN for these 181 weak lines. When we exclude these weak lines, Fig. 5 shows that for the remaining 467 strong lines in the ν_1 fundamental, the CKAPTEN DMS improves on LTP2011 for almost all transitions of Birk *et al.*¹⁹ considered. Lines deemed unstable fall outside of the $\pm 10\%$ deviation shown here and are also excluded; there are 64 unstable theoretical lines. For the transition intensities that LTP2011 underestimates, CKAPTEN fractionally increases

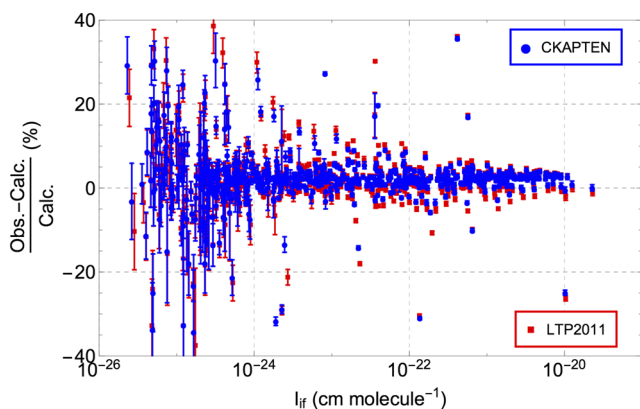


FIG. 4. Log plot of the deviation from measured intensities¹⁹ for all observed ν_1 lines for the CKAPTEN and LTP2011²⁰ DMS. The uncertainties shown are experimental.

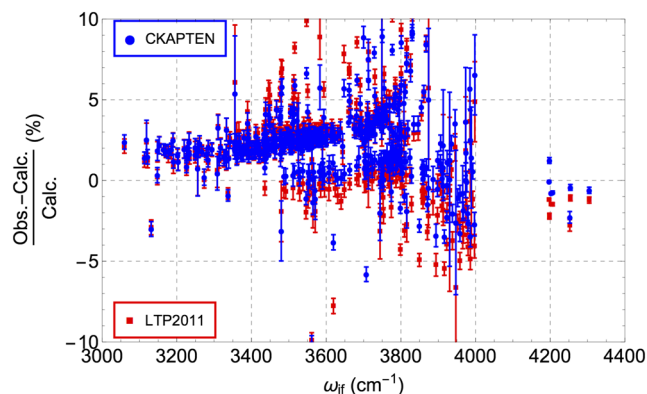


FIG. 5. Comparison of strong ν_1 transitions from CKAPTEN and those of Lodi *et al.*²⁰ against the experimental intensities.¹⁹ The uncertainties are experimental.

these, while it similarly reduces the intensity of those that LTP2011 overestimated, thus highlighting the good stability of CKAPTEN.

CKAPTEN marginally underestimates the same ν_3 fundamental lines that LTP2011²⁰ slightly overestimates (Fig. 6). This overestimation is likely a direct consequence of the larger rms associated with the perpendicular dipole fitting of Lodi *et al.*²⁰

The experiment of Sironneau and Hodges⁵⁰ identified a failure by the BT2 line list²⁹ to accurately model $2\nu_3$ band intensities, with lines overestimated by an average of 5.3%, with those of Schwenke and Partridge (SP2000) averaging less than 1%. The SP2000 line list was actually computed by Tashkun (Tomsk, Russia)⁵² using surfaces due to Schwenke and Partridge.^{23,24} We also find similar problems with the LTP2011 DMS, except that lines are now underestimated by approximately 5% (see Fig. 7).

We compare transitions from CKAPTEN, Lodi *et al.*,²⁰ and the Schwenke and Partridge (SP2000) line list^{23,24,52} with 35 of the experimental⁵⁰ lines in Fig. 7. Of these 35, CKAPTEN is closer to experiment⁵⁰ for 23.

Two outliers at 7745.695 cm^{-1} ($J', K'_a, K'_c = 10\ 7\ 4$) and 7745.866 cm^{-1} ($10\ 7\ 3$) (frequencies and quantum numbers from HITRAN2016) are, respectively, 4% and 2% too weak.

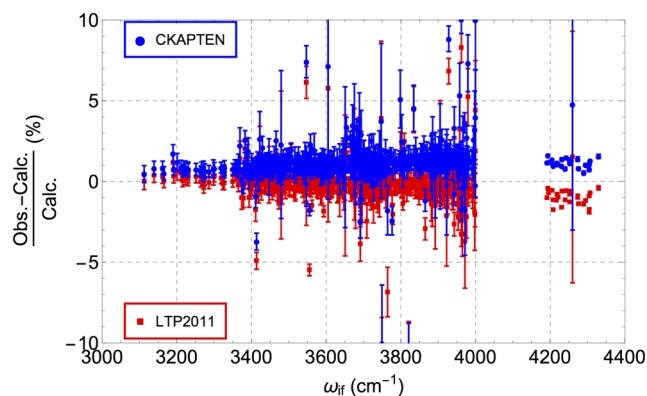


FIG. 6. Comparison of strong ν_3 transitions from CKAPTEN and those of Lodi *et al.*²⁰ against the experimental intensities.¹⁹ The uncertainties are experimental.

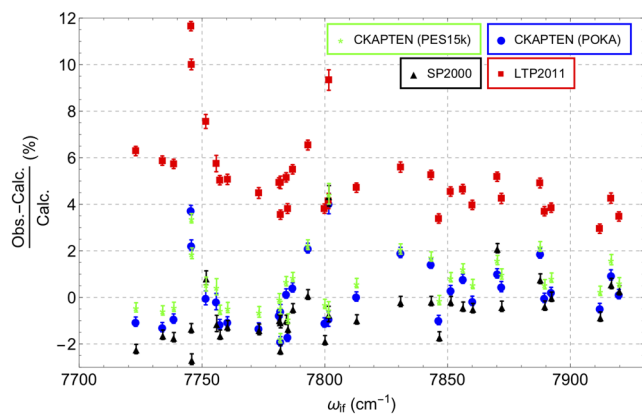


FIG. 7. $2\nu_3$ transitions using the CKAPTEN DMS with both POKAZATEL³⁰ and PES15k⁵¹ potentials, Lodi *et al.*,²⁰ and the Schwenke and Partridge (SP2000) line list^{23,24,52} against the experimental intensities of Sironneau and Hodges.⁵⁰ The uncertainties are experimental.

The transition frequency using the POKAZATEL PES for (10 7 4) is 7745.5 cm^{-1} , while that for (10 7 3) is 7745.674 cm^{-1} , both at a difference of approximately 0.2 cm^{-1} , nearly 7 times the 0.03 cm^{-1} rms reported in Ref. 30. The error in transition frequency is due to the upper energy levels of POKAZATEL, with both lower states reproduced by POKAZATEL³⁰ to $3 \times 10^{-3} \text{ cm}^{-1}$ with respect to the empirical data in HITRAN2016.

Birk *et al.*¹⁹ claim that LTP2011 is accurate to 2% for stable lines, with the notable exception of the (100), (121), (300), and (102) vibrational bands that can deviate by up to 8.8%. CKAPTEN improves the deviation of all these bands, particularly the higher frequency transitions within (121), (300), and (102) (see Table III).

Polyansky *et al.*⁵³ identified a crossing of rotational-vibrational energy levels (020) and (100), which gave rise to a sharing of intensities between transitions involving these states. The trend in the ν_1 fundamental intensity deviations shown in Fig. 5 displays very similar structure to energy level crossings shown in Fig. 8 with a switch in behavior about the crossing region. Given that the LTP2011 DMS²⁰ produces intensities that follow a similar path to ours despite their use

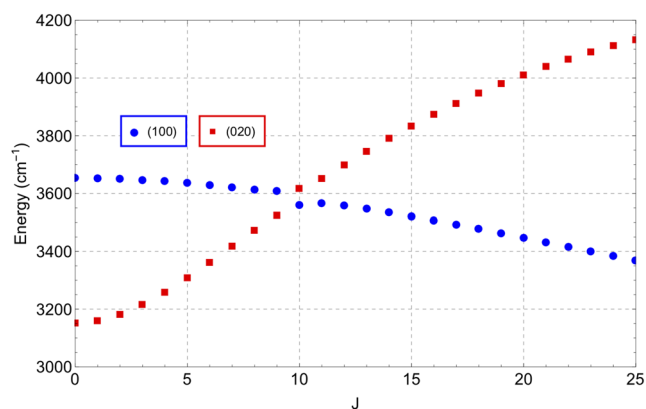


FIG. 8. Energy levels of states (100) and (020) with $J = K_a$ plotted as a function of J . The zero energy is taken as that of the J_0 energy level within the (000) state, as done by Polyansky *et al.*⁵³ Energy levels taken from the MARVEL⁵⁴ database.

of fewer *ab initio* data points, a different functional form, and more parameters, it certainly highlights that the underlying issue is probably with the common potential energy surface.

We tested the latest PES currently available for H_2^{16}O , denoted PES15k, of Mizus *et al.*,⁵¹ which is of high accuracy but only valid below $15\,000 \text{ cm}^{-1}$. Calculations were again constrained to $J \leq 20$ and only transitions lying in the infrared were considered. For the $2\nu_2$ band, we observe that intensities increase in strength and the percentage deviation change by approximately 0.56%, noting that with the POKAZATEL wavefunctions we were approximately 1.85% too weak. For energy levels (10 7 4) and (10 7 3) that are not accurately reproduced by the POKAZATEL PES, PES15k predicts these to approximately 0.03 cm^{-1} .

However, we see the deviation in the ν_1 fundamental go in the opposite direction by 0.40%, with intensities becoming slightly weaker. This behavior again supports our claim of intensity mixing occurring between transitions that involve states ν_1 and $2\nu_2$.

A dramatic improvement is observed in most of the high-energy bands measured by Birk *et al.*¹⁹ PES15k improves upon the POKAZATEL potential for a large number of low-lying bands.

Included in the [supplementary material](#) is a comparison of CKAPTEN and LTP2011²⁰ with all $J = 6-7$ experimentally measured intensities of Birk *et al.*¹⁹

B. Visible and ultraviolet comparisons

Lampel *et al.*¹¹ recorded water vapour absorption in the visible and ultraviolet and compared their observations with several theoretical models; the POKAZATEL line list that utilised the LTP2011S DMS, which is a fewer parameter fit than the larger LTP2011, against the CVR DMS,²⁸ the BT2²⁹ line list, and HITEMP.³¹

For the 450–500 nm window, the strongest line intensities in HITRAN2016 are from the experiment of Tolchenov *et al.*,⁵⁵ with the weaker lines taken from the BT2 line list of Barber *et al.*²⁹ Tolchenov *et al.*⁵⁵ compared their experimental intensities with several other models and determined BT2²⁹ to be the most accurate and reliable in this region.

The two strongest bands in Fig. 9 are (511) and (303), at 487 nm and 471 nm respectively, and theoretical line strengths

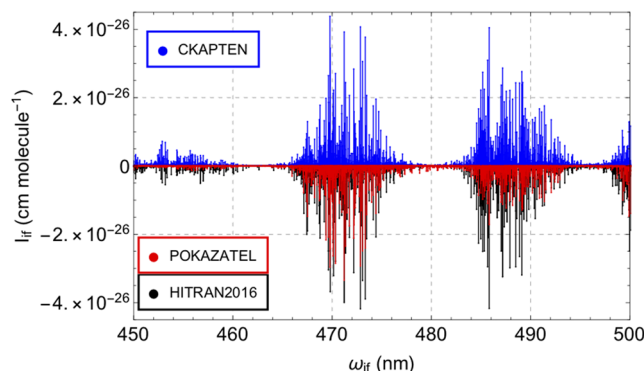


FIG. 9. Intensity comparison of this work, HITRAN2016, and the POKAZATEL line list with transitions constrained to $J_{max} = 14$ for bands (303) and (511).

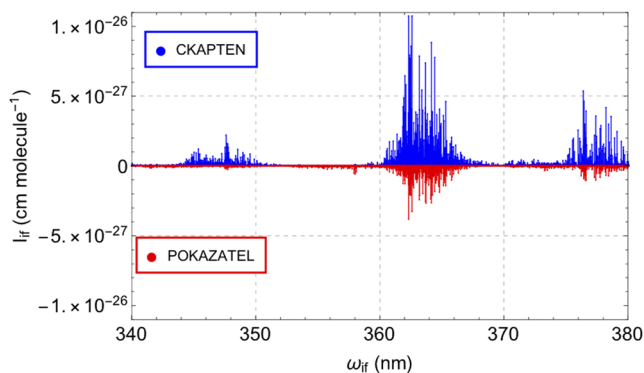


FIG. 10. Comparison of the POKAZATEL line list with the new CKAPTEN surface for 363 nm and 377 nm absorption features.

depend primarily on the accuracy of perpendicular dipole fitting. The strongest line intensities of Tolchenov *et al.*⁵⁵ also share a similar shape to the absorption features measured by Lampel *et al.*¹¹

The POKAZATEL line list fails to accurately model either of the two observed bands in Fig. 9; again, it is likely due to errors in the *ab initio* fitting of LTP2011S perpendicular dipoles, noting a residual of 5×10^{-4} a.u. Comparing CKAPTEN intensities with the strong lines of Tolchenov *et al.*⁵⁵ available in HITRAN2016 which have an associated experimental uncertainty of 2%–5%, the percentage deviation between these bands is in the range of 5%–6%, thus showing good agreement.

Lampel *et al.*¹¹ measured strong water vapour absorption at 363 nm, the (9,0) \pm 0 feature, and determined that it is underestimated in POKAZATEL by a factor of 2.39 ± 0.05 and similarly a weaker band at 377 nm is too small by a factor of 3.1 ± 0.7 .

CKAPTEN predicts the (9,0) \pm 0 structure to be almost 2.7 times larger than that of POKAZATEL (see Fig. 10), approximately 10% greater than the Lampel *et al.*¹¹ observation. Likewise, the magnitude of the 377 nm band is nearly 3.9 times larger than POKAZATEL, slightly outside the error quoted by Lampel *et al.*,¹¹ but given the large uncertainty with their measurement, this is a satisfactory result.

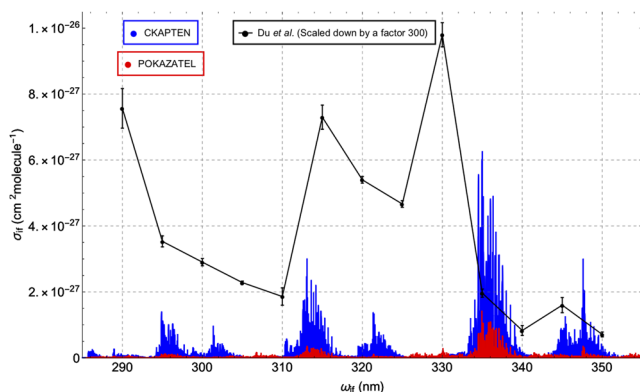


FIG. 11. Cross sections in the 285–355 nm region for the POKAZATEL line list and CKAPTEN surface generated using the HITRAN application program interface (HAPI),⁵⁶ with the scaled observations of Du *et al.*³²

The cross sections measured by Du *et al.*³² have been the subject of much scrutiny^{11,33,57} for being orders-of-magnitude too large. While we also agree that their measurements appear to significantly overestimate the overall absorption in this region, CKAPTEN does, however, predict absorption features that map quite well with their observation (see Fig. 11), with the exception that Du *et al.* show a strong absorption near 290 nm. The band at approximately 335 nm is that with 10 quanta of stretch.

For this band, experimental line positions are blue-shifted by 5 nm compared with theory, while the band at 313 nm is red-shifted by an equal amount. Such large errors remove the POKAZATEL PES as an underlying cause; however, we note that the cross sections of Du *et al.*³² were in fact measured in 5 nm intervals. The POKAZATEL line list also fails to show any significant absorption in this region.

VII. SUMMARY

This new DMS for water vapour, termed CKAPTEN, is the most accurate available and calculations required approximately 80 CPU years to complete. It is based on a grid of 17 628 *ab initio* data points calculated with the Douglas-Kroll-Hess Hamiltonian using the aug-cc-pCV6Z basis set. The computations were calculated at the MR-CI level of theory whose active space included the 9 lowest lying A' molecular orbitals together with two lowest lying A'' molecular orbitals.

Transition intensities that were poorly predicted by previous theoretical models in the infrared region are now closer to experimental measurements than previous theoretical models, notably for bands (002), (300), (102), and (201). For transitions occurring at lower frequencies, theory already predicts intensities to within 2% of the experimental values and CKAPTEN reproduces these intensities with the same or improved accuracy.

High-energy bands, such as (303) and (511), located in the visible spectrum now show excellent agreement with strong experimental lines of Tolchenov *et al.*⁵⁵ available in HITRAN2016, with intensity deviations of about 5%–6%. Both bands also share a similar shape to the observations of Lampel *et al.*¹¹

This new surface predicts the (9,0) \pm 0 feature to within 10% of the atmospheric observation of Lampel *et al.*,¹¹ which is a dramatic improvement over the most recent models that underestimate some bands by a factor of 2.3. The high accuracy associated with CKAPTEN for predicting these intensities will hopefully offer a partial solution to the missing absorber problem in our atmosphere but could only account for possibly a few percentage of extra radiation absorbency.

The crossing of between vibrational energy levels (100) and (020) is identified as the cause of the long-running problem of computing the intensity of ν_1 symmetric stretching fundamental to high accuracy. Use of the latest high accuracy water PES⁵¹ leads to improved results for many transitions.

The stability of the DMS fit is investigated through the removal of half the *ab initio* points. Comparing with accurate experimental measurements in the infrared, we observe

no change in intensities that causes deviation to increase by more than 1% for transition frequencies lower than $10\,000\text{ cm}^{-1}$, up to a maximum change of 1.5% in the (121) band. For observed overtones located in the visible and ultraviolet, we see a maximum change of 12% in the 377 nm band.

Upon testing a new potential surface, we observe intensity changes of approximately 1.35% in the $2\nu_1$ band, 0.56% in the $2\nu_2$ band, 1.40% in $3\nu_1$, and up to 1.63% in $3\nu_3$ transitions. Such changes in these low-lying bands indicate that more work is indeed required in the fitting of potential energy surfaces for the water molecule. A Fortran subroutine containing the CKAPTEN DMS is available as [supplementary material](#) to this article.

Continued validation of the CKAPTEN surface against both experimental and atmospheric data for highly energetic overtones is planned, with the end result being a new line list providing accurate transition intensities extending from the infrared to the near ultraviolet.

SUPPLEMENTARY MATERIAL

The CKAPTEN DMS is available in the [supplementary material](#).

ACKNOWLEDGMENTS

We thank Iouli Gordon and Kelly Chance for helpful discussion during this project and the University College London for their use of both supercomputers, Legion and Grace, which were used for all *ab initio* calculations presented here. This work was partially supported by the UK Natural Environment Research Council under Grant No. NE/N001508/1 and State Project IAP RAS No. 0035-2014-009.

- ¹C. P. McKay, *Proc. Natl. Acad. Sci. U. S. A.* **111**, 12628 (2014).
- ²G. A. Schmidt, R. A. Ruedy, R. L. Miller, and A. A. Lacis, *J. Geophys. Res.* **A 115**, D20106, <https://doi.org/10.1029/2010jd014287> (2010).
- ³M. Wild and B. Liepert, *Environ. Res. Lett.* **5**, 025203 (2010).
- ⁴V. Ramanathan, P. J. Crutzen, J. T. Kiehl, and D. Rosenfeld, *Science* **294**, 2119 (2001).
- ⁵T. G. Huntington, *J. Hydrol.* **319**, 83 (2006).
- ⁶I. E. Gordon, L. S. Rothman, C. Hill, R. V. Kochanov, Y. Tan, P. F. Bernath, M. Birk, V. Boudon, A. Campargue, K. V. Chance *et al.*, *J. Quant. Spectrosc. Radiat. Transfer* **203**, 3 (2017).
- ⁷N. Jacquinet-Husson, R. Armante, N. A. Scott, A. Chédin, L. Crépeau, C. Boutammine, A. Bouhdaoui, C. Crevoisier, V. Capelle, C. Boone *et al.*, *J. Mol. Spectrosc.* **327**, 31 (2016).
- ⁸K. Chance, *Earth Observer* **25**, 10 (2013).
- ⁹S. T. Gulde, M. G. Kolm, D. J. Smith, R. Maurer, G. B. Courrèges-Lacoste, M. Sallusti, and G. Bagnasco, *Proc. SPIE* **10563**, 10563 (2017).
- ¹⁰J. Kim, M. Kim, M. Choi, Y. Park, C.-Y. Chung, L. Chang, S. H. Lee, and G. S. Team, in *IEEE International Geoscience and Remote Sensing Symposium (IGARSS)* (IEEE, Inc., 2016), pp. 4084–4086.
- ¹¹J. Lampel, D. Pöhler, O. L. Polyansky, A. A. Kyuberis, N. F. Zobov, J. Tennyson, L. Lodi, U. Frieß, Y. Wang, S. Beirle *et al.*, *Atmos. Chem. Phys.* **17**, 1271 (2017).
- ¹²Y. Wang, S. Beirle, F. Hendrick, A. Hilboll, J. Jin, A. A. Kyuberis, J. Lampel, A. Li, Y. Luo, L. Lodi *et al.*, *Atmos. Meas. Tech.* **10**, 3719 (2017).
- ¹³T. Wagner, S. Beirle, H. Sihler, and K. Mies, *Atmos. Meas. Tech.* **6**, 2593 (2013).
- ¹⁴H. Wang, X. Liu, K. Chance, G. G. Abad, and C. C. Miller, *Atmos. Meas. Tech.* **7**, 1901 (2014).
- ¹⁵S. Noël, M. Buchwitz, H. Bovensmann, R. Hoogen, and J. P. Burrows, *Geophys. Res. Lett.* **26**, 1841, <https://doi.org/10.1029/1999gl900437> (1999).
- ¹⁶K. Weigel, A. Rozanov, F. Azam, K. Bramstedt, R. Damadeo, K.-U. Eichmann, C. Gebhardt, D. Hurst, M. Kraemer, S. Lossow *et al.*, *Atmos. Meas. Tech.* **9**, 133 (2016).
- ¹⁷S. Noël, S. Mieruch, H. Bovensmann, and J. P. Burrows, *Atmos. Chem. Phys.* **8**, 1519 (2008).
- ¹⁸P. F. Levelt, J. Joiner, J. Tamminen, J. P. Veefkind, P. K. Bhartia, D. C. Stein Zweers, B. N. Duncan, D. G. Streets, H. Eskes, R. van der A *et al.*, *Atmos. Chem. Phys.* **18**, 5699 (2018).
- ¹⁹M. Birk, G. Wagner, J. Loos, L. Lodi, O. L. Polyansky, A. A. Kyuberis, N. F. Zobov, and J. Tennyson, *J. Quant. Spectrosc. Radiat. Transfer* **203**, 88 (2017).
- ²⁰L. Lodi, J. Tennyson, and O. L. Polyansky, *J. Chem. Phys.* **135**, 034113 (2011).
- ²¹L. Lodi and J. Tennyson, *J. Quant. Spectrosc. Radiat. Transfer* **113**, 850 (2012).
- ²²A. A. Kyuberis, N. F. Zobov, O. V. Naumenko, B. A. Voronin, O. L. Polyansky, L. Lodi, A. Liu, S.-M. Hu, and J. Tennyson, *J. Quant. Spectrosc. Radiat. Transfer* **203**, 175 (2017).
- ²³H. Partridge and D. W. Schwenke, *J. Chem. Phys.* **106**, 4618 (1997).
- ²⁴D. W. Schwenke and H. Partridge, *J. Chem. Phys.* **113**, 6592 (2000).
- ²⁵W. Gabriel, E. A. Reinsch, P. Rosmus, S. Carter, and N. C. Handy, *J. Chem. Phys.* **99**, 897 (1993).
- ²⁶U. G. Jørgensen, P. Jensen, G. O. Sørensen, and B. Aringer, *Astron. Astrophys.* **372**, 249 (2001).
- ²⁷J. Lamouroux, S. A. Tashkun, and V. G. Tyuterev, *Chem. Phys. Lett.* **452**, 225 (2008).
- ²⁸L. Lodi, R. N. Tolchenov, J. Tennyson, A. E. Lynas-Gray, S. V. Shirin, N. F. Zobov, O. L. Polyansky, A. G. Császár, J. van Stralen, and L. Visscher, *J. Chem. Phys.* **128**, 044304 (2008).
- ²⁹R. J. Barber, J. Tennyson, G. J. Harris, and R. N. Tolchenov, *Mon. Not. R. Astron. Soc.* **368**, 1087 (2006).
- ³⁰O. L. Polyansky, A. A. Kyuberis, N. F. Zobov, J. Tennyson, S. N. Yurchenko, and L. Lodi, *Mon. Not. R. Astron. Soc.* **480**, 2597 (2018).
- ³¹L. S. Rothman, I. E. Gordon, R. J. Barber, H. Dothe, R. R. Gamache, A. Goldman, V. I. Perevalov, S. A. Tashkun, and J. Tennyson, *J. Quant. Spectrosc. Radiat. Transfer* **111**, 2139 (2010).
- ³²J. Du, L. Huang, Q. Min, and L. Zhu, *Geophys. Res. Lett.* **40**, 4788–4792, <https://doi.org/10.1002/grl.50935> (2013).
- ³³E. M. Wilson, J. C. Wenger, and D. S. Venables, *J. Quant. Spectrosc. Radiat. Transfer* **170**, 194 (2016).
- ³⁴M. Carleer, A. Jenouvrier, A.-C. Vandaele, P. F. Bernath, M. F. Mérianne, R. Colin, N. F. Zobov, O. L. Polyansky, J. Tennyson, and V. A. Savin, *J. Chem. Phys.* **111**, 2444 (1999).
- ³⁵P. Dupre, T. Germain, N. F. Zobov, R. N. Tolchenov, and J. Tennyson, *J. Chem. Phys.* **123**, 154307 (2005).
- ³⁶J. Tennyson, *J. Mol. Spectrosc.* **298**, 1 (2014).
- ³⁷E. S. Medvedev, V. V. Meshkov, A. V. Stolyarov, V. G. Ushakov, and I. E. Gordon, *J. Mol. Spectrosc.* **330**, 36 (2016).
- ³⁸H.-J. Werner, P. J. Knowles, G. Knizia, F. R. Manby, and M. Schütz, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.* **2**, 242 (2012).
- ³⁹L. Lodi and J. Tennyson, *J. Phys. B: At., Mol. Opt. Phys.* **43**, 133001 (2010).
- ⁴⁰T. H. Dunning, Jr., *J. Chem. Phys.* **90**, 1007 (1989).
- ⁴¹D. E. Woon and T. H. Dunning, Jr., *J. Chem. Phys.* **103**, 4572 (1995).
- ⁴²K. A. Peterson and T. H. Dunning, *J. Chem. Phys.* **117**, 10548 (2002).
- ⁴³X. Huang, R. S. Freedman, S. A. Tashkun, D. W. Schwenke, and T. J. Lee, *J. Quant. Spectrosc. Radiat. Transfer* **130**, 134 (2013).
- ⁴⁴A. G. Császár, G. Czako, T. Furtenbacher, J. Tennyson, V. Szalay, S. V. Shirin, N. F. Zobov, and O. L. Polyansky, *J. Chem. Phys.* **122**, 214305 (2005).
- ⁴⁵J. Tennyson, M. A. Kostin, P. Barletta, G. J. Harris, O. L. Polyansky, J. Ramanlal, and N. F. Zobov, *Comput. Phys. Commun.* **163**, 85 (2004).
- ⁴⁶M. Vidler and J. Tennyson, *J. Chem. Phys.* **113**, 9766 (2000).
- ⁴⁷T. Furtenbacher, T. Szidarovszky, J. Hruba, A. A. Kyuberis, N. F. Zobov, O. L. Polyansky, J. Tennyson, and A. G. Császár, *J. Phys. Chem. Ref. Data* **45**, 043104 (2016).
- ⁴⁸I. I. Bubukina, O. L. Polyansky, N. F. Zobov, and S. N. Yurchenko, *Opt. Spectrosc.* **110**, 160 (2011).
- ⁴⁹E. Zak, J. Tennyson, O. L. Polyansky, L. Lodi, S. A. Tashkun, and V. I. Perevalov, *J. Quant. Spectrosc. Radiat. Transfer* **177**, 31 (2016).
- ⁵⁰V. T. Sironneau and J. T. Hodges, *J. Quant. Spectrosc. Radiat. Transfer* **152**, 1 (2015).
- ⁵¹I. I. Mizus, A. A. Kyuberis, N. F. Zobov, V. Y. Makhnev, O. L. Polyansky, and J. Tennyson, *Philos. Trans. R. Soc., A* **376**, 20170149 (2018).

- ⁵²See <http://spectra.iao.ru> for Spectroscopy of atmospheric gases (SPECTRA).
- ⁵³O. L. Polyansky, J. Tennyson, and P. F. Bernath, *J. Mol. Spectrosc.* **186**, 213 (1997).
- ⁵⁴J. Tennyson, P. F. Bernath, L. R. Brown, A. Campargue, M. R. Carleer, A. G. Császár, L. Daumont, R. R. Gamache, J. T. Hodges, O. V. Naumenko *et al.*, *J. Quant. Spectrosc. Radiat. Transfer* **117**, 29 (2013).
- ⁵⁵R. N. Tolchenov, O. Naumenko, N. F. Zobov, S. V. Shirin, O. L. Polyansky, J. Tennyson, M. Carleer, P.-F. Coheur, S. Fally, A. Jenouvrier *et al.*, *J. Mol. Spectrosc.* **233**, 68 (2005).
- ⁵⁶R. V. Kochanov, I. E. Gordon, L. S. Rothman, P. Wcisło, C. Hill, and J. S. Wilzewski, *J. Quant. Spectrosc. Radiat. Transfer* **177**, 15 (2016).
- ⁵⁷J. Lampel, D. Pöhler, J. Tschritter, U. Frieß, and U. Platt, *Atmos. Meas. Tech.* **8**, 4329 (2015).